

# Modeling of Saturated ac Machines Including Air Gap Flux Harmonic Components

Julio C. Moreira, *Member, IEEE*, and Thomas A. Lipo, *Fellow, IEEE*

**Abstract**—A new saturation model for induction machines, which can be easily extended to other types of ac machines, is presented in this paper. It is shown that saturation is responsible for the generation of flux space harmonic components traveling in the air gap with the same synchronous speed as the fundamental flux component with the third being the dominant harmonic component. Superposition of the effects of the fundamental and third harmonic components of the air gap flux is utilized in order to model the saturation of the ferromagnetic parts of the machine. The concept of winding functions is used to derive the inductance terms relating to both stator and rotor winding components. In this approach, the air gap length is assumed to be variable, being a function of the position and level of the air gap flux. Terminal and torque values for steady and transient states are obtainable from the proposed model, with experimental results showing that the model proposed predicts spatial saturation effects with good accuracy.

## I. INTRODUCTION

SATURATION phenomenon in ac electrical machines, especially in induction machines, has received considerable attention in the past few years. In many cases, modeling of ac machines under saturation is based on a small signal linearization around an operating point [1]. Another approach proposed in [2] utilizes the concept of reorienting the  $d$ - $q$  axis in order to incorporate the effects of spatially dependent saturation into the main flux path. Although the effects of main flux modeling based on the fundamental component of MMF has received much attention, the generation of higher flux harmonics due to saturation has not received the same attention in the literature. The problem of saturation harmonics was first described in [3], wherein the main concerns of the author were the prediction of the air gap flux harmonic component amplitudes and the changes in machine steady state performance. Another approach to determine the air gap flux density harmonic components for induction machines under saturated conditions is described in [4].

In this paper, a new saturation model for ac machines is developed based on the fact that the saturation of the magnetic field paths is responsible for introducing harmonic components of flux due to the nonlinear nature of the saturation phenomenon. A  $d$ - $q$  model is derived from the conven-

tional constant parameter ac machine model, which is modified to account for the saturation. This modification consists of making the air gap length a function of the air gap flux position and amplitude. It is shown that as a consequence of the saturation, a third harmonic flux component exists, and third harmonic currents are induced in the rotor circuit. Although they are of low amplitude, they contribute to the total torque and introduce rotor losses. In addition, the third harmonic of flux when linking the stator windings produces a zero sequence third harmonic voltage component in the stator phases. When the three-phase voltages are summed, the fundamental and characteristic harmonics are canceled and the resultant waveform contains, mainly, a third harmonic voltage signal, which can be used as a means of locating the machine air gap flux [5], [6].

A detailed analytical model is presented, and a digital simulation program is introduced. The resulting model can be used to predict the dynamic behavior of the machine during transients with large signal variations. Experimental results show that the model proposed predicts the saturation effects accurately.

## II. INCORPORATING SATURATION IN THE CONSTANT PARAMETER MODEL

To commence the analysis, it is useful to consider a wye-connected three-phase ac machine when supplied from a sinusoidal voltage supply. In this case, the machine will clearly not have zero sequence harmonic components in the stator current. The air gap flux, on the other hand, will have a harmonic content that is a function of the machine saturation. Two different saturation effects must then be distinguished according to the place where the saturation occurs. When saturation occurs in the stator and rotor teeth, the resultant air gap flux assumes a flattened sinusoidal form with a peak value  $B_g$ , as indicated in Fig. 1.

The resultant air gap flux in Fig. 1 has a harmonic content that includes all odd harmonic components, including the triplens third, ninth, and so forth. A Fourier series expansion of the air gap flux, however, indicates that the third harmonic is the dominant harmonic component at all saturation levels. Hence, saturation of the teeth paths can be modeled by the superposition of a third harmonic to the fundamental component of the air gap field flux. It is important to note that this third harmonic flux component travels in the air gap with the same velocity and direction as the fundamental air gap flux component, always keeping the phase synchronism with the fundamental.

Saturation can also occur at the stator and rotor cores, also

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J. C. Moreira is with Whirlpool R&E, Benton Harbor, MI 49022.

T. A. Lipo is with the Department of Electrical and Computer Engineering, University of Wisconsin, Madison, WI 53706-1691.

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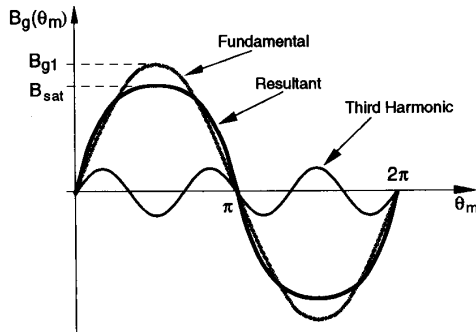


Fig. 1. Air gap flux density distribution for a stator/rotor teeth saturation condition (heavy line). The fundamental and third harmonic components are also shown.

causing the introduction of harmonic components in the air gap flux. In this case, the core flux distribution assumes a flattened sinusoidal form, whereas the air gap flux distribution assumes a peaked waveform. This peaked air gap flux distribution waveform is explained by the fact that the air gap flux density  $B_g(\theta_m)$  is proportional to the space derivative of the core flux density  $B_c(\theta_m)$ . Again, a series of odd harmonics will be present in the air gap flux distribution, but the third harmonic is once again the dominant harmonic component. In this case, the third harmonic has a phase relationship with the fundamental such that one of its peaks enforces the fundamental, thus giving the peaked shape for the resultant air gap flux.

Saturation of the stator and rotor teeth are more common in practical machines since the quantity of iron in the core is generally greater than in the teeth, where a much higher flux density exists. Hence, the saturation condition exemplified by Fig. 1 is more typical and will be the case assumed in the analysis throughout this paper.

From the previous paragraphs, one can conclude that the presence of the third harmonic component in the air gap flux is a sufficient but not a necessary condition for the occurrence of saturation. Both saturation mechanisms have an opposite effect as far as third harmonic flux component is concerned. Saturation of the core sections will produce a third harmonic flux component, which is in direct phase opposition to the third harmonic produced by the saturation of the teeth paths. Consequently, a highly saturated machine does not necessarily have a high third harmonic flux content.

The third harmonic air gap flux linking the stator spatial winding third harmonic component is responsible for the induction in the stator phases of a zero sequence third harmonic voltage. Since no third harmonic currents will flow in a wye-connected machine, this third harmonic voltage keeps the same phase relationship with the flux harmonic and, hence, can be utilized to locate the air gap field flux with respect to any of the machine terminal quantities. It is important to note that the third harmonic air gap flux component keeps the same zero phase relationship with the fundamental flux for any load condition. This is a direct consequence of the fact that when saturation occurs, the air gap flux keeps its flattened sinusoidal form independently on the

machine load. It can also be argued in another way by saying that the third harmonic rotor currents are not large enough to produce a rotor mmf that would change the position of the third harmonic flux with respect to the fundamental flux component. This result is of capital importance when control applications, as proposed in [5] and [6], utilizing the third harmonic voltage signal are implemented.

Reduction of the magnetic permeability of the iron paths is the immediate consequence of saturation. Fig. 2 shows, in a simplified manner, the variation of the stator and rotor iron permeability for a two-pole machine. Saturation of the stator and rotor teeth is assumed to occur only in the iron parts with the flux density defined by the resultant component of the air gap flux  $B_g$ . Core saturation regions are also shown in Fig. 2, but for the analysis proposed here, only the teeth saturation will be taken into account as explained before. As the resultant flux  $B_g$  moves around the air gap, so does the spatial variation of the permeability, suggesting that the permeability variation can be viewed as a traveling wave in the air gap.

We can also recognize that when teeth saturation occurs, the decrease in permeability of the iron paths has the effect of increasing the teeth reluctances for the region around the resultant component of the air gap flux. Keeping the necessary assumption of constant and infinite iron permeability, the reluctance variation of the teeth can be viewed as an air gap reluctance variation. Since the air gap has a constant permeability  $\mu_0$ , the gap length has to be made variable to accommodate the reluctance variation due to the saturation. Therefore, instead of having a variation in the iron permeability, we then assume a pseudo air gap length variation as the flux saturates the iron parts.

From the discussion in the preceding paragraph, it becomes clear that saturation of the stator and rotor teeth can be incorporated into the linear model for the general rotating machine as described in [7] by making the air gap length a function of the saturation level and spatial position. Although the air gap length variation is not necessarily a sinusoidal function, it will be assumed so in this analysis. A typical air gap length variation for a pole pitch is shown in Fig. 3, together with its inverse function  $g^{-1}(\theta_m)$ . The inverse gap function  $g^{-1}(\theta_m)$  is defined for convenience in dealing with multiplication of functions rather than division, as it would be the case in using  $g(\theta_m)$  for computing the inductances of the machine. Note that  $g_o$  corresponds to the nominal machine air gap length.

As depicted in Fig. 3, an inverse gap function can be defined as

$$g^{-1}(\theta) = k_e + k_m \cos(2\theta_m - 2\theta_f) \quad (1)$$

where

$$k_e = g_o^{-1} - k_m$$

with  $\theta_m$  and  $\theta_f$  representing, respectively, an angle measure around the air gap periphery and the position of the air gap flux with respect to the phase  $a$  reference axis.

The factor  $k_m$  is derived from the saturation factor for the

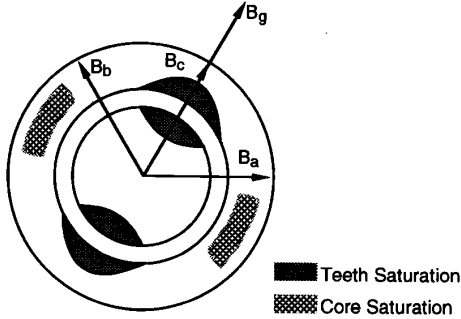


Fig. 2. Cross sectional view of a saturated machine showing the variation of the iron permeability in the stator and rotor parts. Flux density components and the total air gap flux density  $B_g$  are also shown.

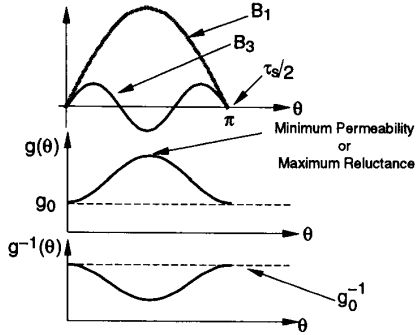


Fig. 3. Air gap length variation function  $g(\theta_m)$  and its inverse  $g^{-1}(\theta_m)$ .

machine  $k_{sat}$  as

$$k_m = \frac{2 k_{sat} - 1}{3 k_{sat} g_0} \quad (2)$$

with  $k_{sat}$  being the ratio between the fundamental components of the air gap voltage for a nonsaturated and saturated conditions.

The self and mutual inductances of the stator winding are now computed based on the fact that the air gap length is variable. Based on this approach, winding functions for the stator windings are described as

$$N_{as}(\theta_m) = N_{s1} \cos \theta_m + N_{s3} \cos 3\theta_m \quad (3)$$

$$N_{bs}(\theta_m) = N_{s1} \cos(\theta_m - 2\pi/3) + N_{s3} \cos 3\theta_m \quad (4)$$

$$N_{cs}(\theta_m) = N_{s1} \cos(\theta_m + 2\pi/3) + N_{s3} \cos 3\theta_m. \quad (5)$$

The self inductance for one of the stator phases, phase  $a$  for instance, is computed as

$$L_a = \frac{\lambda_{aa}}{i_a} = \mu_0 r l \int_0^{2\pi} g^{-1}(\theta_m) [N_{as}(\theta_m)]^2 d\theta_m \quad (6)$$

with  $r$  and  $l$  representing the average air gap radius and the lamination stack length, respectively.

The mutual inductance between windings of phases  $a$  and  $b$ , for instance, is computed from

$$L_{ab} = \frac{\lambda_{ab}}{i_b} = \mu_0 r l \int_0^{2\pi} g^{-1}(\theta_m) N_{as}(\theta_m) N_{bs}(\theta_m) d\theta_m. \quad (7)$$

The response of the rotor cage to the two components of the air gap flux is a cage current that has two rotating components: one at fundamental and the other at three times the fundamental frequency. Hence, a polyphase fundamental and third harmonic current components are induced in the rotor, creating two components of rotor mmf. It is then necessary to assume that the rotor contains two sets of polyphase winding functions as

$$N_{ar1} = N_{r1} \cos(\theta_m - \theta_r) \quad (8)$$

$$N_{br1} = N_{r1} \cos(\theta_m - \theta_r - 2\pi/3) \quad (9)$$

$$N_{cr1} = N_{r1} \cos(\theta_m - \theta_r + 2\pi/3) \quad (10)$$

$$N_{ar3} = N_{r3} \cos(3\theta_m - 3\theta_r) \quad (11)$$

$$N_{br3} = N_{r3} \cos(3\theta_m - 3\theta_r - 2\pi/3) \quad (12)$$

$$N_{cr3} = N_{r3} \cos(3\theta_m - 3\theta_r + 2\pi/3). \quad (13)$$

Computation of the stator inductances as described above yields, for a three-phase machine, three self and six mutual inductances for the stator windings, which are collected in a matrix defined as  $\bar{L}_s$ . The rotor self and mutual inductances are also computed utilizing the winding functions as describe above. These inductance values are organized in two inductance matrixes: one for the fundamental rotor winding function  $\bar{L}_{r1}$  and other for the third harmonic  $\bar{L}_{r3}$ . These inductances will be time and rotor speed dependent and will be a function of the saturation factor. The constants  $N_{s1}$ ,  $N_{r1}$  represent, respectively, the number of turns for the stator and rotor circuits, whereas  $N_{r3}$  is defined as the equivalent number of turns for the rotor third harmonic circuit.

The next step in the modeling of the machine is the derivation of the mutual inductances between the stator and rotor windings. This derivation process leads to two inductance matrixes ( $\bar{L}_{rs1}$  and  $\bar{L}_{rs3}$ ) containing the mutual coupling elements between fundamental and third harmonic components of the rotor and stator windings. These mutual inductances are also time and saturation level dependent.

The  $abc$  phase variable machine equations in the vector form are then derived as follows:

$$\bar{v}_{abc} = r_s \bar{i}_{abc} + \frac{d\lambda_{abc}}{dt} \quad (14)$$

$$\bar{v}_{abc r1} = r_{r1} \bar{i}_{abc r1} + \frac{d\lambda_{abc r1}}{dt} \quad (15)$$

$$\bar{v}_{abc r3} = r_{r3} \bar{i}_{abc r3} + \frac{d\lambda_{abc r3}}{dt}. \quad (16)$$

The flux linkage equations are written in terms of the stator and rotor current components and the self and mutual inductance matrixes previously defined:

$$\bar{\lambda}_{abc} = \bar{L}_s \bar{i}_{abc} + \bar{L}_{sr1} \bar{i}_{abc r1} + \bar{L}_{sr3} \bar{i}_{abc r3} \quad (17)$$

$$\bar{\lambda}_{abc r1} = \bar{L}_{r1} \bar{i}_{abc r1} + \bar{L}_{rs1} \bar{i}_{abc} \quad (18)$$

$$\bar{\lambda}_{abc r3} = \bar{L}_{r3} \bar{i}_{abc r3} + \bar{L}_{rs3} \bar{i}_{abc}. \quad (19)$$

Note from the flux linkage expressions above that there is

no mutual coupling between rotor fundamental and third harmonic windings, as expected.

### III. SATURATION MODEL IN D-Q-N VARIABLES

The solution for the model derived in Section II requires the knowledge of the stator and rotor voltage vectors  $\bar{v}_{abc}$ ,  $\bar{v}_{abc r1}$ , and  $\bar{v}_{abc r3}$  since they are the independent variables for the model. The third harmonic component present in these voltages, however, precludes the utilization of them as independent variables because they will be a function of the saturation level of the machine. To circumvent this problem, it is necessary to transform the machine equations to the  $dqn$  general reference frame. In this new reference system, the phase voltages are free from the unknown third harmonic voltage component because these components form a zero sequence set that is eliminated from the  $d-q$  components of the phase voltages. On the other hand, the  $n$  component of the stator phase voltage is not zero as it is usually the case in the conventional machine model; it contains the third harmonic voltage component, which will be solved from the  $n$ -axis component of the stator flux. To maintain generality, the transformation is carried out in a general reference frame rotating at a speed given by  $\omega$ .

The machine model in the  $dqn$  general reference frame is given by

$$v_{qs} = r_s i_{qs} + \frac{d\lambda_{qs}}{dt} + \omega \lambda_{ds} \quad (20)$$

$$v_{ds} = r_s i_{ds} + \frac{d\lambda_{ds}}{dt} - \omega \lambda_{qs} \quad (21)$$

$$v_{ns} = r_s i_{ns} + \frac{d\lambda_{ns}}{dt} \quad (22)$$

$$v'_{qr1} = r'_{r1} i'_{qr1} + \frac{d\lambda_{qr1}}{dt} + (\omega - \omega_r) \lambda_{dr1} \quad (23)$$

$$v'_{dr1} = r'_{r1} i'_{dr1} + \frac{d\lambda_{dr1}}{dt} - (\omega - \omega_r) \lambda_{qr1} \quad (24)$$

$$v'_{nr1} = r'_{r1} i'_{nr1} + \frac{d\lambda_{nr1}}{dt} \quad (25)$$

$$v'_{qr3} = r'_{r3} i'_{qr3} + \frac{d\lambda_{qr3}}{dt} + 3(\omega - \omega_r) \lambda_{dr3} \quad (26)$$

$$v'_{dr3} = r'_{r3} i'_{dr3} + \frac{d\lambda_{dr3}}{dt} - 3(\omega - \omega_r) \lambda_{qr3} \quad (27)$$

$$v'_{nr3} = r'_{r3} i'_{nr3} + \frac{d\lambda_{nr3}}{dt} \quad (28)$$

where the rotor-primed variables are referred to the stator by a stator/rotor turns ratio transformation. For a wye-connected stator, no zero sequence currents will circulate, and consequently,  $i_{ns}$  is zero, and the same can be assumed for the rotor currents  $i_{nr1}$  and  $i_{nr3}$ .

The flux linkages in the  $dqn$  stator frame are derived after a long algebraic process such as

$$\lambda_{qs} = \left[ L_{ls} + L_{m1} + \frac{L_{m1} k_m}{2k_e} \cos 2(\theta_f - \theta) \right] i_{qs}$$

$$\begin{aligned} & - \frac{L_{m1} k_m}{2k_e} \sin 2(\theta_f - \theta) i_{ds} \\ & + \left[ L_{m1} + \frac{L_{m1} k_m}{2k_e} \cos 2(\theta_f - \theta) \right] i'_{qr1} \\ & - \frac{L_{m1} k_m}{2k_e} \sin 2(\theta_f - \theta) i'_{dr1} \\ & + L_{s13} \cos 2(\theta_f - \theta) i'_{qr3} + L_{s13} \sin 2(\theta_f - \theta) i'_{dr3} \end{aligned} \quad (29)$$

$$\begin{aligned} \lambda_{ds} = & \left[ L_{ls} + L_{m1} - \frac{L_{m1} k_m}{2k_e} \cos 2(\theta_f - \theta) \right] i_{ds} \\ & - \frac{L_{m1} k_m}{2k_e} \sin 2(\theta_f - \theta) i_{qs} \\ & + \left[ L_{m1} - \frac{L_{m1} k_m}{2k_e} \cos 2(\theta_f - \theta) \right] i'_{dr1} \\ & - \frac{L_{m1} k_m}{2k_e} \sin 2(\theta_f - \theta) i'_{qr1} \\ & + L_{s13} \cos 2(\theta_f - \theta) i'_{dr3} + L_{s13} \sin 2(\theta_f - \theta) i'_{qr3} \end{aligned} \quad (30)$$

$$\begin{aligned} \lambda_{ns} = & \sqrt{2} L_{s13} \cos(2\theta_f + \theta) i_{qs} + \sqrt{2} L_{s13} \sin(2\theta_f + \theta) i_{ds} \\ & + \sqrt{2} L_{s13} \cos(2\theta_f + \theta) i'_{qr1} \\ & + \sqrt{2} L_{s13} \sin(2\theta_f + \theta) i'_{dr1} \\ & + \sqrt{2} L_{m3} \cos 3\theta i'_{qr3} + \sqrt{2} L_{m3} \sin 3\theta i'_{dr3}. \end{aligned} \quad (31)$$

Similar flux linkage equations are obtained for the two rotor circuits. It is clear from (29) to (31) that the flux linkages for the  $d$ - and  $q$ -axis components do not result in decoupling between the  $d-q$  axis variables as in the conventional model. The machine saturation is responsible for the loss of the decoupling between variables as well as for the generation of the zero sequence flux component  $\lambda_{ns}$ . The model is greatly simplified by making  $\theta = \theta_f$  when the general rotating  $d-q$  axes rotate synchronously with the air gap flux, as shown by (32) to (34). By doing so, the cross coupling between the  $d-q$  variables observed in the equations above is eliminated. On the other hand, the zero sequence flux linkage still depends on all stator and rotor current components:

$$\begin{aligned} \lambda'_{qs} = & (L_{ls} + L_{m1} + L_{ksat}) i'^f_{qs} \\ & + (L_{m1} + L_{ksat}) i'^f_{qr1} + L_{ksat} i'^f_{qr3} \end{aligned} \quad (32)$$

$$\begin{aligned} \lambda'_{ds} = & (L_{ls} + L_{m1} - L_{ksat}) i'^f_{ds} \\ & + (L_{m1} - L_{ksat}) i'^f_{dr1} + L_{ksat} i'^f_{dr3} \end{aligned} \quad (33)$$

$$\begin{aligned} \lambda'_{ns} = & \sqrt{2} L_{s13} i'^f_{qs} \cos 3\theta_f + \sqrt{2} L_{s13} i'^f_{ds} \sin 3\theta_f \\ & + \sqrt{2} L_{s13} i'^f_{qr1} \cos 3\theta_f + \sqrt{2} L_{s13} i'^f_{dr1} \sin 3\theta_f \\ & + \sqrt{2} L_{m13} i'^f_{qr3} \cos 3\theta_f + \sqrt{2} L_{m13} i'^f_{dr3} \sin 3\theta_f. \end{aligned} \quad (34)$$

The superscript  $f$  indicates that the  $dqn$  reference frame

rotates synchronously with the air gap flux, and the inductance terms are defined as follows:

$$L_{m1} = \frac{3}{2} L_{ms} = \frac{3}{2} k_l k_e N_{s1}^2 \quad (35)$$

$$L_{ksat} = \frac{L_{m1}}{2} \frac{k_m}{k_e} \quad (36)$$

$$L_{m13} = \frac{3}{2} k_l k_e N_{s1} N_{s3} \quad (37)$$

$$L_{s13} = \frac{3}{4} L_{s3} = \frac{3}{4} k_l k_m N_{s1} N_{s3} \quad (38)$$

with the factor  $k_l$  defined as

$$k_l = \mu_0 r l \pi.$$

An equivalent circuit for the induction machine operating under saturation is developed from the model equations above. Such a circuit is depicted in Fig. 4 for the case where the rotating reference frame is fixed in the air gap flux. It incorporates a stator and two rotor circuits for each one of the  $d$  and  $q$  variables. Note that the zero sequence voltage  $v_{ns}^f$  is simply obtained as the derivative of the flux  $\lambda_{ns}^f$  since no zero sequence currents flow in the circuit.

As in [2], the resulting model is dependent on the air gap flux position defined as  $\theta_f$ . Consequently, the  $d$ - $q$ - $n$  variable transformation is performed with respect to this position, when the general argument  $\theta$  becomes  $\theta_f$ . A similar technique as in [2] to orient the rotating  $d$ - $q$  frame with respect to the air gap flux is used in the implementation of the simulation model.

The torque produced by the machine is computed from the mutual inductance terms and stator and rotor currents as in [7]. Two torque components arise from this computation process. One of these components is related to the interaction between the fundamental stator and rotor current components, whereas the other torque term is related to the interaction between the fundamental stator and third harmonic rotor currents. The expressions for these two torque components are given by (39) and (40):

$$T_{e1} = \left( \frac{P}{2} \right) \left( \frac{3}{2} \right) (L_{m1} - L_{ksat}) (i_{qs}^f i_{dr1}^f - i_{ds}^f i_{qr1}^f) \quad (39)$$

$$T_{e3} = \left( \frac{P}{2} \right) \left( \frac{9}{2} \right) L_{s13} (i_{qs}^f i_{dr3}^f - i_{ds}^f i_{qr3}^f). \quad (40)$$

The effects of saturation on the torque production mechanism is clearly shown by the torque expressions above. As expected, as the saturation level increases, the torque due to the interaction of the fundamental stator and rotor current components decreases. It is interesting to note that saturation is also responsible for the induction of third harmonic currents that will produce a third harmonic rotor flux component. These harmonic rotor currents interact with the stator fundamental current via the pseudo air gap length variation giving rise to the torque component in (40), which is normally very small compared with the torque produced by the

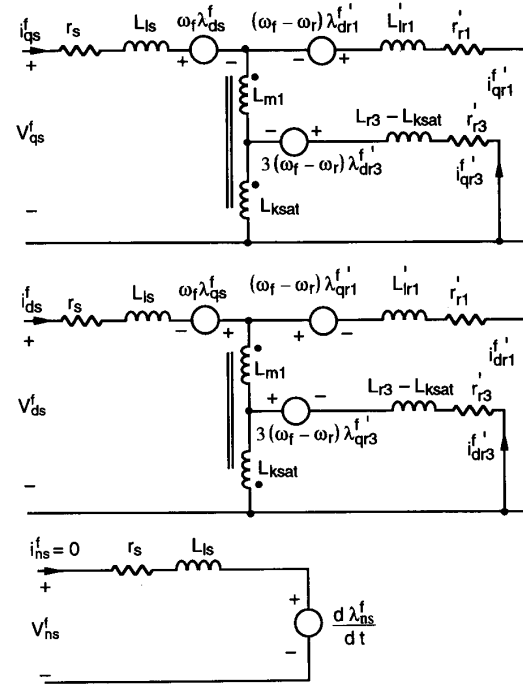


Fig. 4. Equivalent circuit for a saturated induction machine.

fundamental components. These harmonic rotor currents are also responsible for extra copper losses in the rotor.

The inductance term  $L_{m1}$  corresponds to the unsaturated value of the magnetizing inductance, which is experimentally obtained. The other inductance terms are given in terms of the machine design parameters.

#### IV. EXPERIMENTAL AND SIMULATION RESULTS

A 3-hp, 230-V, four-pole induction machine has been instrumented in order to evaluate the correlation between experimental and theoretical results from the model described. All the experimental tests and simulation results were obtained with the machine excited from a sinusoidal voltage supply.

The machine electrical parameters were obtained via the conventional no-load and locked rotor tests. Some of the machine design parameters, like number of turns for the stator windings and the number of rotor bars, are needed for evaluation of the inductance coefficients used in the model.

The no-load characteristic curve (stator phase rms voltage versus stator rms current) is measured in the laboratory and is plotted in Fig. 5. The no-load characteristic curve obtained from the saturation model is plotted in the same figure. It is clear that the model predicts the saturation with very good accuracy. In general, the proposed model predicts less magnetizing current, which is an indication that higher harmonic components of the air gap flux contribute to the saturation of the machine. The inclusion of these higher harmonic components of flux will considerably increase the complexity of the saturation model with very little improvement in the final result.

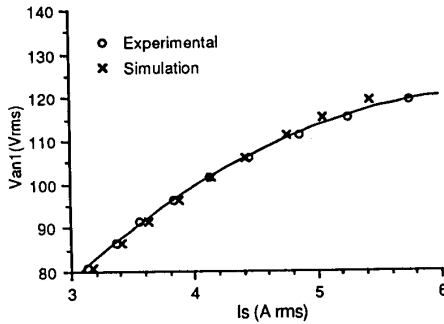


Fig. 5. Magnetization characteristic curve for the induction motor measured experimentally and as predicted by the saturation model.

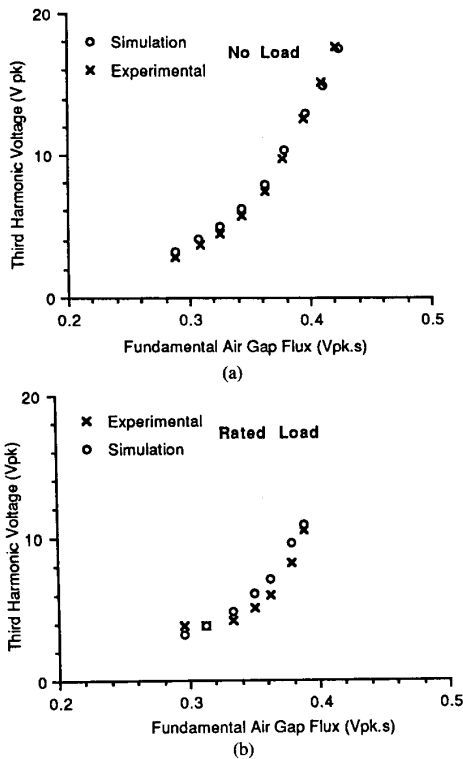


Fig. 6. Amplitude of the stator third harmonic voltage signal as a function of the amplitude of the air gap flux fundamental component for different motor mechanical load conditions. Simulation and experimental results are shown for (a) no-load and (b) rated-load conditions.

Fig. 6 shows the experimental and simulation results for the amplitude of the stator third harmonic voltage signal as a function of the amplitude of the air gap flux fundamental component for different motor mechanical load conditions. Again, the good correlation between experimental and theoretical results shows that the saturation model predicts with accuracy the third harmonic component generated by saturation. Note that the third harmonic voltage signal is of substantial magnitude even for flux levels around 60% of the rated value. Consequently, the third harmonic voltage signal will be present even when the machine operates in a field-weakening condition.

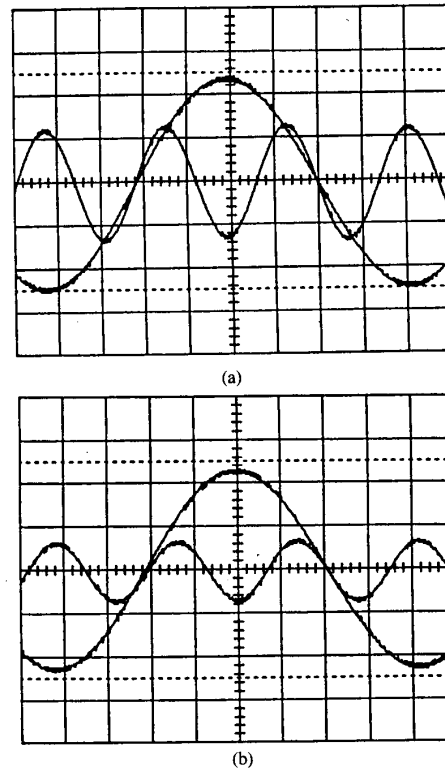


Fig. 7. Fundamental and third harmonic components of the air gap flux for conditions of (a) no-load and (b) rated-load. Scales: fundamental flux— $0.2 \text{ V} \cdot \text{s}/\text{div}$ ; third harmonic flux— $5 \text{ mV} \cdot \text{s}/\text{div}$ ; time— $2 \text{ ms}/\text{div}$ .

Fig. 7 shows the fundamental and the third harmonic components of air gap flux measured by a search coil installed in the air gap of the machine. These plots are obtained for no-load and rated load conditions at rated voltage. Note that, as predicted by the model, the loading of the machine does not interfere with the phase relationship between the two flux components. This comes from the fact that the third harmonic rotor currents are very small so that the rotor third harmonic mmf is not large enough to change the flattened shape of the air gap flux. Simulation results have shown that the rotor third harmonic currents have values lower than 1% of the rated stator current.

Finally, Fig. 8 depicts the frequency spectrum of the air gap flux density, measured by means of a search coil introduced in the air gap, for the induction machine running at no-load condition. As predicted, the third harmonic is larger than the other harmonics, thus justifying the assumption made that it is the dominating harmonic component.

## V. CONCLUSION

A new saturation model for an induction machine, including the effect of third harmonic spatial saturation, has been developed and presented in this paper. The model is derived from a conventional constant parameter ac machine model modified to account for the stator and rotor teeth saturation effects. It can be used to predict the dynamic behavior of the machine for large signal variations. The model is based on

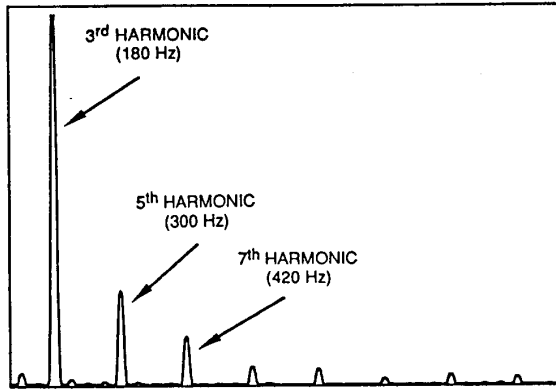


Fig. 8. Frequency spectrum for the air gap flux showing that the third harmonic is the dominant component. Scales: 6 mV · s full scale, 1100 Hz full scale.

the fact that saturation introduces harmonic components in the air gap flux with the third being the dominant harmonic component. A pseudo variation of the air gap length is assumed to account for the modulation of the fundamental air gap flux by the third harmonic component. The resultant model contains time-varying inductance coefficients that are also a function of the saturation level represented by the saturation factor  $k_{sat}$ . It is shown that the rotating third harmonic air gap flux component is responsible for the induction of a zero-sequence third harmonic stator phase voltage that can be utilized as a means to locate the air gap flux, as described in [5] and [6].

APPENDIX A

Table I shows the induction motor parameters used in this work.

TABLE I  
INDUCTION MOTOR PARAMETERS

Quantity	Symbol	Value
Line voltage	V1	230 V rms
Output power	P <sub>0</sub>	3.0 hp
Speed	w <sub>r</sub>	1740 r/min
Poles	P	4
Stator resistance	r <sub>s</sub>	1.11 Ω
Rotor resistance	r <sub>r</sub>	0.47 Ω
Stator leakage reactance	X <sub>ls</sub>	1.05 Ω
Rotor leakage reactance	X <sub>lr</sub>	1.05 Ω
Unsaturated magnetic reactance	X <sub>m</sub>	22.09 Ω
Rotor inertia	J <sub>m</sub>	0.0104 Kg-m <sup>2</sup>
Number of rotor slots	n <sub>r</sub>	46
Number of stator slots	n <sub>s</sub>	36
Rotor skew	—	1 slot
Stator pole pitch	τ <sub>s</sub>	7/9
Load inertia	J <sub>L</sub>	0.0200 Kg-m <sup>2</sup>
Airgap length	g <sub>0</sub>	0.4165 mm
Rotor stack length	l <sub>m</sub>	65.4 mm

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**Julio C. Moreira** (S'81-M'90) was born in Sao Paulo, Brazil. He received the B.E.E. and M.E.E. degrees in 1979 and 1983, respectively, from the State University of Campinas (UNICAMP). He received the Ph.D. degree in electrical engineering from the University of Wisconsin, Madison, in 1990.

He worked as an Engineer and later as a Vice Coordinator of the Electric Drive Systems Research Laboratory at UNICAMP from 1980 to 1985. In 1981, he became a Researcher and an Assistant Professor at that same institution, teaching courses in electrical machines, power electronics, and controls. His research activities included the development of dc and ac motor drives for an electric vehicle. He also served as a consultant in the areas of electric drive systems and switched power supplies. He later joined the Whirlpool Corporate Research and Engineering Center, Benton Harbor, MI, where he is involved with research in the areas of power electronics, electric drive systems, controls, and motor design.

Dr. Moreira is a member of ETA, NOΣ, IAS, PES, and IES.



**Thomas A. Lipo** (M'64-SM'71-F'87) received the B.E.E. and M.S.E.E. degrees from Marquette University, Milwaukee, WI, in 1962 and 1964, respectively, and the Ph.D. degree in electrical engineering from the University of Wisconsin in 1968. He was an NRC Postdoctoral Fellow at the University of Manchester Institute of Science and Technology, Manchester, England, from 1968 to 1969.

From 1969 to 1979, he was an Electrical Engineer in the Power Electronics Laboratory of Corporate Research and Development of the General Electric Company, Schenectady, NY. He became Professor of Electrical Engineering at Purdue University, Lafayette, IN, in 1979 and later joined the University of Wisconsin, Madison, in the same capacity. He has been involved in the research of power electronics and ac drives for over 25 years.

Dr. Lipo has received 11 IEEE prize paper awards including corecipient of the Best Paper Award in IEEE TRANSACTIONS ON INDUSTRY APPLICATIONS for 1984. In 1986, he received the Outstanding Achievement Award from the IEEE Industry Applications Society for his contributions to the field of ac drives.